Twisted Bi-Layer Graphene: Microscopic Rainbows


Twisted bi-layer graphene (tBLG) has recently attracted interest due to the peculiar electrical properties that arise from its random rotational configurations. Our experiments on CVD-grown graphene from Cu foil and transferred onto Si substrates, with an oxide layer of 100 nm, reveal naturally-produced bi-layer graphene patches which present different colorations when shined with white light. In particular yellow-, pink- and blue-colored areas are evidenced. Combining optical microscopy, Raman spectroscopy and transmission electron microscopy we have been able to assign these colorations to ranges of rotational angles between the two graphene layers. Optical contrast simulations have been carried out, proving that the observation of the different colorations is due to the angle-dependent electronic properties of tBLG combined with the reflection that results from the layered structure tBLG/100 nm-thick SiO$_2$/Si. Our results could lead the way to an easy selective identification of bi-layer graphene merely through the observation on an optical microscope.

Carbon atoms covalently bonded and arranged in a honeycomb lattice constitute graphene. When two of these layers are stacked on top of each other a new material takes the scene, namely bi-layer graphene. Although bi-layer and multi-layer graphene have been studied along with single-layer graphene since its observation in 2004,[1] the early reports consisted on mechanically exfoliated bi-layer, which maintained the AB stacking of HOPG.[2,3] AB stacking or Bernal stacking is considered as the most stable configuration is obtained by overlapping two graphene sheets in such a way that a carbon atom in one layer would sit on the center of the hexagonal ring from the other layer, corresponding to a rotational angle of 60°. However, when two graphene layers are randomly placed on top of each other, the overlap can generate a new superlattice in the form of a Moiré pattern. Different twist angles between the layers will generate different Moiré patterns which can be visualized using high resolution transmission electron microscopy (HRTEM)[4] or scanning tunneling microscopy (STM).[5,6] In particular, Raman spectroscopy plays a very important role in identifying the stacking nature of graphene. In 2006, Ferrari and co-authors[2] demonstrated that bi- and few-layer graphene derived from HOPG by mechanical exfoliation have a unique Raman fingerprint. It was later that scientists realized that artificially produced, chemical vapor deposition (CVD) grown and epitaxially grown bi-layer graphene samples had different stacking orders and hence different Raman spectra.[7-11] It has been proven that the Moiré superlattices give rise to periodic potentials that alter the electronic dispersion of the structure. According to recent publications the extent to which the electronic properties are modified depends on the twist angle. For large angles the layers are decoupled and behave like single layer graphene (θ > 15°), which is the case for few-layer graphene grown on SiC.[12] For intermediate angles (2° < θ < 15°) the band velocity is decreased and van Hove singularities appear in the density of states (DOS), where the energy gap varies as a function of θ. For very small angles (θ < 2°) the Fermi velocity becomes angle dependent and sharp peaks appear close to the Dirac point[13,14] Those van Hove singularities have been predicted theoretically and observed experimentally through scanning tunneling spectroscopy (STS)[15-18] and angle-resolved photoemission spectroscopy (ARPES).[19]

The optical properties of single layer graphene, Bernal-stacked and twisted bi-layer graphene are governed by their electronic structure. While single layer and Bernal-stacked bi-layer graphene show a constant absorption in the visible range (~2.3% and ~4.6%, respectively),[20] for tBLG, the many van Hove singularities that appear in the DOS will alter accordingly the light absorption behavior for each rotational angle.[21]

Producing artificial bi-layer graphene consists on intentionally placing two layers of graphene on top of each other, for instance by mechanically folding graphene[22,23] or by transferring two CVD single-layer graphene sheets on a target substrate.[24] Carozo and co-workers[23] showed that by dragging graphene with an AFM tip it can be folded onto itself, artificially producing bi-layer graphene with well defined rotational angles. They reported on the superlattices thus formed and the appearance of new Raman spectroscopy features in the vicinity of the G-band, dependent on the angle of rotation, namely R and R’ bands. Kim and co-workers[24] artificially produced it by superimposing two CVD-grown graphene layers. Since CVD-grown graphene has a polycrystalline nature, they were able to produce bi-layer graphene.
with a wide variety of rotational angles within the same sample. Coupling transmission electron microscopy (TEM) studies and Raman spectroscopy they showed striking differences in the electronic structure of twisted bi-layer graphene (tBLG).

The chemical vapor deposition of large area graphene using metals as catalysts produces mono-layer graphene and, inevitably, bi-layer patches.\textsuperscript{[8–10]} It has been proven that these bi-layer patches present differences in rotational angles due to the polycrystalline nature of the grown graphene.\textsuperscript{[25–27]} Havener and co-authors\textsuperscript{[27]} showed, using dark-field transmission electron microscopy and Raman spectroscopy, that areas of bi-layer CVD-grown graphene can be composed of domains with different rotational angles. They demonstrated an angle-dependent resonance of the G-band and showed that this resonance is also energy-dependent by probing the sample with different laser excitation energies. The above mentioned studies\textsuperscript{[22–24,27]} demonstrate that Raman spectroscopy constitutes a mighty tool to probe the angles of rotation in tBLG.

In this work we prove that bi-layer graphene produced by low pressure CVD transferred to SiO\textsubscript{2}/Si substrates with a 100 nm oxide layer, shows differences in the optical reflection when observed with an optical microscope. Using Raman spectroscopy we have been able to find G-band resonances in these regions, which indicate that the color differences are directly related to the rotational angles between layers. This conclusion was further confirmed with TEM studies, and supported by simulations of the optical reflectance contrast for different $\theta$ values of tBLG.

The inspection of our samples by optical microscopy revealed peculiar colorations in the two layer patches that naturally appear in CVD-grown graphene. Figure 1 shows micrographs of our samples where different colorations are visible in particular zones of the bi-layer graphene islands. Although the colorations can be observed without any treatment of the images, for the sake of clarity we have enhanced the contrast using software tools to render them more obvious (untreated images can be found in Figure S1 of the Supporting Information).

Our investigations reveal that the thickness of the silicon oxide layer is of paramount importance for the observation of these colorations. Besides our observations on Si substrates with oxide thicknesses of 100 nm, we have transferred our samples to Si substrates with oxide thicknesses of 300 nm and 90 nm. For the 300 nm SiO\textsubscript{2} thickness, no colorations were found, not even after software treatment of the images. Subtle increases in contrast were found for some regions of the bi-layer patches where the G-band turned out to be resonant with $E_{\text{laser}}$ = 2.41 eV ($\lambda$ = 514 nm, Figure S2 C and D), as previously discussed in Reference [27]. Optical micrographs and Raman spectra of samples transferred to Si substrates with 300 nm oxide layer can be found in Figure S2 of the Supporting Information. In the case of 90 nm, pink colorations are the most evident and yellow- and blue-colored areas can be also identified but show a decrease in contrast (results not included here).

Although the differences in absorption properties are intrinsic to tBLGs dictated by their unique angle-dependent DOS, the colorations are not observed unless the optimal SiO\textsubscript{2} thickness is used. The observation of the different colorations is due to the angle-dependent electronic properties of tBLG combined with the reflection that results from the layered structure tBLG/100 nm-thick SiO\textsubscript{2}/Si.

The reflectance of single- and multi-layer graphene on SiO\textsubscript{2} substrates has been widely studied\textsuperscript{[20,26,29]} however, to the best of our knowledge, this is the first time that different colorations are observed in bi-layer graphene. The colorations observed can mainly be catalogued as yellowish, bluish, and pinkish. We have performed a large number of Raman spectroscopy measurements of these particular areas using three different laser excitation energies. Figure 2A shows a representative micrograph of our samples where the above mentioned areas (yellow, blue, pink) are observed. Numbers 1, 2 and 3 point out the spots where measurements were recorded. Our measurements reveal that yellow areas show a strong G-band resonance with $E_{\text{laser}}$ = 2.54 eV ($\lambda$ = 488 nm, spot 1 in Figure 2A), blue areas show a G-band resonance when probed with $E_{\text{laser}}$ = 1.96 eV ($\lambda$ = 633 nm, spot 2 in Figure 2A) and for pink areas the G-band resonance is found with both, $E_{\text{laser}}$ = 2.41 eV ($\lambda$ = 514 nm) and $E_{\text{laser}}$ = 2.54 eV ($\lambda$ = 488 nm) (spot 3 in Figure 2A).

We have checked these color-dependent and energy-dependent resonances of the G-band at a larger scale by recording Raman maps. Figure 2B illustrates the area where maps were acquired, the image to the right corresponds to the same image after software treatment, where zones of interest with pink, yellow and blue colorations are visible. The corresponding G-band intensity Raman maps are shown below the optical images, contours of the bi-layer areas have been included as guides to the eye. A careful cross-inspection of optical images and Raman maps reveals that pink and yellowish areas show the highest intensities when using $E_{\text{laser}}$ = 2.41 eV ($\lambda$ = 514 nm) and that the intensity is enhanced in bluish areas when using $E_{\text{laser}}$ = 1.96 eV ($\lambda$ = 633 nm).

To investigate the structural nature of these particularly-colored patches, we have performed TEM experiments. A
A detailed inspection of our Raman spectra in the G-band region, allowed us to locate the angle-dependent R-bands, previously reported by Carozo and coworkers.\[23\] Different measurements in the blue zones, casted R-bands with different frequencies (see Figure 3C for representative curves drawn in red), suggesting a window of rotational angles for this particular coloration of tBLG. The same phenomena were observed for pink and yellow zones (see Figure 3C green and blue curves, respectively, for representative spectra). We realized that the position of the R-bands is color-dependent, as depicted in the plot from Figure 3C, where the frequency regions corresponding to each color have been highlighted. Calculating the correspondence of the R-bands to rotational angles\[30\] we estimated the angles that can be assigned to our different colored bi-layer graphene patches.

Angles of 9° to 11° correspond to bluish areas, angles of 11° to 13° correspond to pinkish areas and angles of 13° to 15° correspond to yellowish areas. The above angle assignment matches perfectly our TEM results, where an angle of 12.8° was found for a pink zone and an angle of 13.4° for a yellow zone.

Our observations are validated as well by comparison with results recently published by K. Kim et al.\[24\] and R.W. Havener et al.\[27\]. In the above mentioned reports, the G-band resonance was found for $E_{\text{\text{\text{\text{laser}}} = 1.96 \text{eV}}}$ at $\sim 10^\circ$, such angle of rotation corresponds to a bluish zone where we, as well, found G-band resonance for $E_{\text{\text{\text{\text{laser}}} = 1.96 \text{eV}}}$ for $E_{\text{\text{\text{\text{laser}}} = 2.41/2.33 \text{eV}}}$ their resonance was found at $\sim 13^\circ$, corresponding to a pink zone, where we find G band resonance for $E_{\text{\text{\text{\text{laser}}} = 2.41 \text{eV}}}$ (see Figure 2).

We have further confirmed our observations by simulations of the optical contrast of bi-layer graphene using a model that considers the reflection of a system air/bi-layer graphene/SiO$_2$/Si. Details of the model can be found in the Supporting Information. In order to qualitatively reproduce the effect of the van Hove singularities (vHS) in the optical absorption of tBLG, we have calculated the dielectric function using the graphene refractive index and incorporated a causal Gaussian absorption profile using a Kramers-Kronig analysis. We have estimated the position of the absorption band for each rotation angle by calculating the $\theta$-dependent energy separation of the vHS in tBLG using an empirical formula derived by I. Brihuega et al.\[37\]
graphene samples can yield selective angle-dependent fabrication of devices and thus enable the tunable control of electrical and optical properties.

**Experimental Section**

Graphene samples were produced by CVD at low pressures with methane as carbon source and copper as catalyst. The produced samples were transferred onto SiO$_2$/Si (oxide thickness 100 nm) substrates to be able to optically identify single and bi-layer graphene and to perform Raman measurements.

Raman spectroscopy measurements were performed on a LabRam Horiba instrument with laser excitation energies of 2.54 eV, 2.41 eV and 1.96 eV. Powers have been kept below 1 mW and low and high resolution gratings have been used in our measurements (150 and 1800 g/mm, respectively).

To carry out the TEM studies, the SiO$_2$/Si transferred samples have been re-transferred to Quantifoil grids via potassium hydroxide.$^{[31]}$ The sample was annealed at 200 °C for 10 min before introducing into the TEM. HRTEM and selected-area electron diffraction experiments were performed using an aberration-corrected FEI Titan 80–300 operated at 80 kV. The HRTEM images were obtained at Scherzer conditions with a spherical aberration value of 0.02 mm.

**Supporting Information**

Supporting Information is available from the Wiley Online Library or from the author.

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