Metallic bonds remain one of the most important and least understood of the chemical bonds. In this study, we generated Re$_2$ molecules in which the Re–Re core is unsupported by ligands. Real-time imaging of the atomic-scale dynamics of Re$_2$ adsorbed on a graphitic lattice allows direct measurement of Re–Re bond lengths for individual molecules that changes in discrete steps correlating with bond order from one to four. Direct imaging of the Re–Re bond breaking process reveals a new bonding state with the bond order less than one and a high-amplitude vibrational stretch, preceding the bond dissociation. The methodology, based on aberration-corrected transmission electron microscopy imaging, is shown to be a powerful analytical tool for the investigation of dynamics of metallic bonding at the atomic level.

**RESULTS**

Dirhenium molecules are generated from Re$_2$(CO)$_{10}$ compound inserted into SWNTs, followed by thermal or $\epsilon^-$-beam–induced elimination of CO ligands (Fig. 1A). The dissociation of Re$_2$(CO)$_{10}$ molecules and the removal of CO ligands is confirmed by Fourier transform infrared spectroscopy (FTIR) measurements (fig. S1). If several precursor molecules Re$_2$(CO)$_{10}$ are agglomerated in the nanotube, they form a metal nanocluster with several Re atoms under these conditions (16, 17); in the case of a single isolated molecule, Re$_2$(CO)$_{10}$ ligand elimination yields a well-defined Re$_2$ molecule as observed by aberration-corrected high-resolution TEM (AC-HRTEM) (three typical images of Re$_2$ are shown in Fig. 1A). Two individual dirhenium molecules can be observed in the same SWNT in the top AC-HRTEM image of Fig. 1A. Monitoring an individual Re$_2$ by time-series 80-kV AC-HRTEM imaging (Fig. 1B) shows that the two Re atoms are bonded to each other within the dirhenium molecule, as Re$_2$ moves as a whole in the cavity of SWNT, which differs notably from the dynamics of individual heavy atoms in graphitic structures that move by hopping from one vacancy defect in the carbon lattice to another under $\epsilon$-beam irradiation (18–21). The Re–Re bond ensures that the movement of both atoms is correlated, and the dirhenium molecule takes positions in the SWNT, consistent with small number of possible orientations: (i) chemisorbed standing (20 s), (ii) chemisorbed lying (0 s), and (iii) successive chemisorbed standing and lying configurations during the exposure time (26 s) (Fig. 1B). The same Re$_2$ molecule is able to change its position and orientation over the duration of the experiment. This process is stimulated by the $\epsilon$-beam serving as both a source of energy and the imaging probe, as the kinetic energy of the incident electrons is transferred to the Re atoms (directly $\sim$1.0 eV and via carbon atom $\sim$14.7 eV). As the positions of the rhenium atoms are monitored continually during the exposure time of each frame (1 s), the motion of Re$_2$ switching from one orientation to another can be captured within a single frame and analyzed (see Fig. 2). For example, it is possible to discern that the change in the bonding orientation that takes place during the 26th second corresponds to a change from

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the initial standing orientation lasting 0.29 s into the lying orientation lasting 0.71 s in the 1.0 s frame (fig. S2).

**DISCUSSION**

There is a well-known strong inverse correlation between the M–M bond length and bond order (3). The Re–Re single σ bond length in Re₂(CO)₁₀ is 0.304 nm for the free molecule and 0.302 nm in the crystal (22). However, in our experiments, after the CO ligands have been removed, the Re–Re bond measures 0.220 nm (± 0.015 nm) at the initial state (0 s; Fig. 1B), indicating that the bond order is close to four, as it agrees well with the length of a quadruple Re–Re bond in [Re₂X₈]²⁻ of 0.224 nm (2). Our density functional theory (DFT) calculations for Re₂ match well with the experimentally observed distances, predicting Re–Re bond lengths of 0.21 nm for a free (nonbonded) Re₂ and approximately 0.22 nm for the dirhenium molecule bound to the SWNT wall by the stable standing orientation. Because the SWNTs are cylindrical surfaces, all the bond lengths quoted in this work are corrected from the directly measured projected distance using the method described in fig. S3. Statistical analysis of Re–Re bond lengths in 226 frames from 0 to 553 s (Fig. 2C) shows three major values of the length of 0.22, 0.25, and 0.30 nm, occurring with 44, 24, and 11% probability (Fig. 2E). The nonrandom, discrete values assumed by the Re–Re bond are likely to be correlated with the quadruple bond, double bond, and single bond, respectively, which were previously reported from bulk measurements on various complexes containing a Re₂ core (3, 22). In additional 18% of the cases, the bond length cannot be measured as marked “N/A,” as the “standing” state is aligned with the direction of the electron beam. Being limited by the spatial resolution of only spherical AC-TEM, the atomic structure of the SWNT and hence the precise position of the Re atoms with respect to the carbon lattice remain unresolved.

With the help of the Cc/Cs-corrected SALVE TEM, the e⁻-beam–stimulated dynamics of dirhenium molecule are recorded with a much higher spatial resolution, as shown in movie S2. A Re₂ molecule can be observed standing at defective sites on the top SWNT of two parallel SWNTs at the initial state (0 s; Fig. 2A shows the first stage of movie S2 from 0 to 87 s). After 50 s, Re₂ shows similar behavior to the Re₂ in Fig. 1B, switching between multiple different orientations at the defective site. An intact SWNT is impermeable to almost all atoms and ions except H⁺ (23). However, a vacancy defect on the SWNT could provide a channel for atoms to penetrate. At 50 s, one Re atom of the dirhenium molecule goes across the upper wall of the host SWNT from the defective site. In the following 20 s, the whole dirhenium molecule climbs out and starts drifting along the outer surface of the host SWNT, driven by stronger bonding of the metal with defect from outside, as was predicted by dynamic simulations.
From 72 to 85 s, the dirhenium molecule drifts along the outer surface of the host SWNT and then slides into the one-dimensional vdW gap, as indicated by yellow arrows (Fig. 2B). The simulated structure in the right of Fig. 2B shows an atomic structure of the lying orientation of the dirhenium molecule when it is drifting along the outer surface of host SWNT. The position of the Re atoms are indicated by yellow arrows. The AC-HRTEM images are partially covered by the model of an SWNT with chirality of (10, 10). The simulated structure shows the position of dirhenium molecule in the frame of 82 s. (C) Corrected bond length changing of the right dirhenium molecule in movie S1 (Titan Re₂ A); error is ±0.015 nm. (D) Corrected bond length of the dirhenium molecule drifting along the outer wall of SWNT in movie S2 (SALVE Re₂); error is ±0.010 nm. (E) The histogram shows the count frequencies of the different bond lengths of Titan Re₂ A (0 to 553 s).

The vdW gap between the two parallel SWNTs in the initial state of the second stage (89 s in Fig. 3A) is approximately 0.32 nm (± 0.010 nm), consistent with the theoretical prediction by Lennard-Jones model (25). This vdW gap confines Re₂ and provides a one-dimensional channel in the vacuum, limiting Re₂ translation and rotation. In the first 7 s (87 to 94 s) after the dirhenium molecule moves into the vdW gap, the contrast of the Re atoms shows unusual symmetrically elongated features, indicating that the Re atoms are delocalized on the time frame of a single exposure (0.5 s), which may be caused by oscillations along the Re–Re bond, described as a stretch vibration (Fig. 3B). The center distance of the elongated spots is 0.36 nm, which is slightly longer than the Re–Re single bond (around 0.30 nm) in compounds (21), while the amplitude of the oscillation of Re atoms is ±0.04 nm. After oscillating for 7 s, Re₂ breaks into two unbound Re atoms with a distance of 0.58 nm (94 s in Fig. 3A and the second column in Fig. 3B). After the separation, the contrast of two Re atoms is no longer elongated and becomes near round. This observation indicates that once the Re–Re bond length exceeds that of a single bond, the molecule may exist in an excited state manifested in a large amplitude of vibrations, and that beyond this critical point, the bond breaks and the two metal atoms separate. In our experiments, Re₂ trapped in the one-dimensional vdW gap acts like a harmonic oscillator, pumped with the kinetic energy transferred from the 80-keV electrons to the Re atom, until it is large enough to overcome the barrier of Re–Re bond dissociation (at 94 s in Fig. 3A). The two Re atoms move freely in the one-dimensional
vdW gap until they collide and recombine into Re₂ at 97 s. The bond distance of the reformed Re₂ reduces from 0.30 (97 s) to 0.23 nm (132 s). The contrast of two Re atoms in the frame of 132 s are near round with a distance of 0.23 nm, indicating that the two Re atoms are bound again by a quadruple bond. This reformed dirhenium molecule shows the same properties as before, moving between different orientations in the one-dimensional vdW gap in the following 34 s, similar to the behavior inside and on the SWNT.

In addition, the free Re atom is chemically active toward SWNT, as shown in our previous work (26), and interacts with the upper SWNT restructuring the carbon lattice at 97 s. At 200 s, Re₂ dissociates again and recombines at 212 s. This dirhenium molecule moves in the vdW gap in the following 181 s and merges into the large Re cluster in the lower SWNT as shown in movie S2.

DFT calculations have been carried out to elucidate the stable configurations and electronic structure of Re₂ bonded to SWNT (Fig. 4). As discussed in Materials and Methods, a search of possible stable bonding configurations inside the SWNT revealed only the presence of a single standing mode with a Re–Re bond length of 0.22 nm, in which Re₂ bonds in an η⁶-mode to all six carbon atoms of a single hexagon with a “4 + 2” asymmetry introduced by the curvature of the nanotube wall. The binding energy of this configuration is −165 kJ/mol (−1.71 eV), relative to the separated dirhenium molecule and SWNT. On the outside of the nanotube, the η⁶-standing configuration is again found to be the lowest energy orientation, but the change in curvature introduces several stable lying configurations, which are local minima at higher energy, geometrically and energetically similar to one another (fig. S7).

To understand the bonding configurations observed, it is instructive to examine the spatial extent of the molecular orbitals involved in the bonding. The atomic Re orbitals and the occupied and low-lying unoccupied orbitals of Re₂ are presented in fig. S4, and the electronic structure of the dirhenium molecule is discussed in more detail in the Supplementary Materials. The highest occupied molecular orbital (HOMO) [singly occupied molecular orbital (SOMO)] of the bound species in the standing configuration corresponds to the HOMO...
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Fig. 4. Interactions of Re$_2$ with carbon nanotube. (A) Comparison of the spatial extent of the Re$_2$ HOMO (SOMO) viewed end-on and side-on; the overlap with electron density of the SWNT LUMO sufficient for bonding is only achieved when Re$_2$ approaches the nanotube wall in the end-on orientation. (B) Side-on and end-on views of the HOMO (SOMO) of the bound configuration, showing the contraction of the (blue) orbital lobes that maximizes bonding overlap (in the end-on view, some carbon atoms have been removed and the electron density isovalue has been increased to aid visualization). (C) Illustration of the dynamics of the dirhenium molecule in the SWNT, with the associated calculated energies relative to the separated molecule and SWNT (left), and front and side views of each state.

(SOMO) of the isolated Re$_2$, but with contraction of the orbital lobes that are perpendicular to the nanotube axis (Fig. 4B). This maximizes the overlap with the two carbon atoms that are closer to the dirhenium due to the curvature-induced 4 + 2 asymmetry (shown in blue in Fig. 4A). The “end-on” standing η$^6$-bonding is more stable than “side-on” lying bonding, as this provides a better orbital overlap between the Re$_2$ HOMO and the SWNT LUMO, in terms of size, shape, and ease of the above orbital contraction. On the outside of the SWNT, the rhenium dimer encounters the convex curvature of the nanotube, and two carbon atoms lie further from the dimer than the other four. The orbital contraction is therefore less pronounced, and the larger orbital lobes presented by the side-on orientation of dirhenium are able to take part in stable bonding with the nanotube wall. Note that the curvature of the SWNT wall is, in turn, affected by the dirhenium molecule, exhibiting a slight flattening of the wall in the region of the bonding, as can be seen in the side views in Fig. 4C. Our calculations predict the presence of standing configurations inside the nanotube; the presence of stable lying configurations could be caused by the different curvature of SWNTs with different diameters. As the bonding modes available to Re$_2$ seem to be largely influenced by the curvature of the SWNT wall, it is reasonable to assume that the wider nanotube with a smaller degree of curvature imposes less of a restriction on the possible bonding modes and behaves more similarly to the outside of the tube. In addition, the presence of defects and deformations in the SWNT under electron irradiation were not considered in our simulations but may affect the bonding between dirhenium and the SWNT.

Using our experimental observations in conjunction with the theoretical modeling, Fig. 4C illustrates a plausible energy profile for the movement of Re$_2$ along the nanotube. We propose that, upon elimination of the CO ligands from the precursor molecule, Re$_2$ bonds to the interior of the SWNT wall, a process we predict to be strongly thermodynamically (~165 kJ/mol binding energy) and kinetically (no energy barrier) favorable. On the order of seconds in our conditions, the electron beam transfers energy to excite Re$_2$, from which the molecule can relax into adjacent energy minima (Fig. 4C). The absence of any local minima in which Re$_2$ exists inside the SWNT without being bonded to the walls explains the observed stepwise motion of the dirhenium molecule under electron irradiation, as these would result in a more long-range migration process.

We show that the simultaneous function of the electron beam as a source of energy and an imaging tool allows advancement in the understanding of metallic bonding. The rhenium homonuclear diatomic molecules Re$_2$ are created in nanotubes in situ under AC-HRTEM observation, and dynamics of Re$_2$ are imaged with atomic resolution. The metallic bond can change over time as a result of the electron beam effect, most frequently adopting a quadruple bond order. Re$_2$ can reversibly dissociate to independent atoms under...
the electron beam, which transfers up to \(-1\) eV (\(-100\) kJ/mol) to a Re atom, directly observed when the molecule is confined in the one-dimensional gap between two parallel NTs. An unexpected bonding state where the Re-Re bond order is less than one, accompanied with a large amplitude of vibration observed just before the bond dissociation, indicates the existence of previously unknown bonding states between metal atoms. A combination of DFT calculations and AC-HRTEM time-resolved experiments sheds light on the atomistic mechanism of metallic bond dynamics in real time and direct space, opening up avenues for discovery of new inorganic and organometallic species by electron microscopy methods, such as predicted W\(_5\) with quintuple bond or Re\(_5\) clusters, and exploration of their physical properties and chemical reactivity at the level of the single atom.

**MATERIALS AND METHODS**

**Synthesis of dirhenium molecules confined in SWNTs**

SWNTs (Carbon Solutions, USA) were thermally treated to open the ends of the nanotubes and to remove the residual amorphous carbon from the outer walls of the SWNTs before use. The as-received SWNTs were heated in air for 30 min at 600°C with a weight loss of approximately 30% observed for samples. The freshly opened SWNTs were mixed with a twofold by weight excess of Re\(_2\)(CO)\(_{10}\) compound (Sigma-Aldrich). The resultant mixture was then sealed in a Pyrex glass tube under reduced pressure (10\(^{-6}\) mbar) and heated to 150°C for 3 days. The sample was then rapidly cooled to room temperature, and the SWNTs were washed with tetrahydrofuran (100 ml) to remove any species from the outside of the nanotube and filtered through a polytetrafluoroethylene membrane (pore size = 0.2 μm). The sample was then dispersed in isopropanol and drop cast onto a lacy carbon-coated copper TEM grid for HRTEM analysis. The sample was then dispersed in isopropanol and drop cast onto a lacy carbon-coated copper TEM grid for HRTEM analysis. The formation of dirhenium molecules can also be achieved by heating Re\(_2\)(CO)\(_{10}\)@SWNT at high temperature. This was achieved by sealing Re\(_2\)(CO)\(_{10}\)@SWNT under argon in a Pyrex tube and heating at 500°C for 3 hours. Rhenium cluster@SWNT samples prepared in this way were identical to the samples formed via electron beam–induced elimination of CO ligands from Re\(_2\)(CO)\(_{10}\). FTIR measurement shown in fig. S1 confirms the removal of CO from nanotubes during the heating process.

**Electron microscopy**

AC-HRTEM images in movie S1 were carried out using an image side C\(_{60}\)-corrected FEI Titan 80-300 TEM operated at 80 kV accelerating voltage with a modified filament extraction voltage for information limit enhancement. Images were recorded on a slow-scan charge-coupled device camera type Gatan UltraScan 1000XP. Because a certain amount of electron dose was required to achieve atomic resolution, the exposure time for each frame was 1.0 s. The dose rate of e-beam applied in movie S1 was 1.05 × 10\(^8\) e\(^-\) nm\(^{-2}\) s\(^{-1}\).

AC-HRTEM images in movie S2 were carried out using the image side C\(_{60}/C_\text{Si}\)-corrected SALVE instrument (20 to 80 kV) operated at 80 kV accelerating voltage with a point resolution of <0.08 nm. The exposure time for the frames in movie S2 was 0.5 s. The dose rate of e-beam applied in movie S1 was 1.05 × 10\(^8\) e\(^-\) nm\(^{-2}\) s\(^{-1}\). The dose rate of e-beam applied in movie S2 was 1.10 × 10\(^8\) e\(^-\) nm\(^{-2}\) s\(^{-1}\).

For each observation, approximately 30 s were spent for adjusting the magnification, dose rate, and focal length. All imaging experiments were carried out at room temperature. In addition, 26 s (from 252 to 278 s) was spent for centering the drifted SWNT in movie S1 during continuous irradiation.

**Computational methods**

DFT simulations were performed to find stable bonding configurations of the dirhenium molecule with the host SWNT. Rhenium atoms were modeled with the BLYP/SRSC functional/effecctive core potential pairing, with carbon and hydrogen atoms modeled with a BLYP/6-31G\(^*\) functional/basis set pair for initial geometry optimizations, followed by optimizations using BLYP/6-311G\(^*\) (see the Supplementary Materials for more information on the choice of DFT method). Grimme’s empirical DFT-D3(BJ) model (27) was used to account for the effects of dispersion. The reported binding energy was corrected for basis set superposition error by using the counterpoise method. The hydrogenated (8,8) SWNT shown in fig. S6 was used as a model nanotube, with a diameter of 1.1 nm compared to the experimental (16,1) SWNT with a diameter of 1.3 nm. The optimized rhenum dimer was placed in initial positions corresponding to those in fig. S5 as well as several different nonbonded positions in the center of the SWNT, and all atom positions (including Re\(_2\)) were fully optimized. On the inside of the SWNT, all optimizations resulted in the formation of a standing bonding configuration centered over a hexagon in the tube (shown in Fig. 4). On the outside of the tube, the equivalent hexagon-centered standing bonding configuration was the lowest energy structure, but there were also several additional bonding configurations (shown in fig. S7). The Mulliken charge distribution of the bonded standing state found in the DFT calculations is shown in fig. S8.

**REFERENCES AND NOTES**


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Imaging an unsupported metal–metal bond in dirhenium molecules at the atomic scale
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