

9. Exercise sheet for the course “Modeling with PDEs”
 (Hydrodynamic and drift-diffusion models, related equations)

Exercise 1. The following system of PDEs is called the *bipolar model*:

$$\partial_t n + q^{-1} \operatorname{div}_x J_n = -R(n, p), \quad J_n = -q\mu_n(U\nabla_x n + n\nabla_x V), \quad (1)$$

$$\partial_t p - q^{-1} \operatorname{div}_x J_p = -R(n, p), \quad J_p = q\mu_p(U\nabla_x p - p\nabla_x V), \quad (2)$$

$$R(n, p) = \frac{np - n_i^2}{\tau_p(n + n_d) - \tau_n(p + p_d)}, \quad (3)$$

$$-\varepsilon_s \Delta V = q(n - p - C). \quad (4)$$

Here n, p are the electron and hole densities (resp.), q is the electron charge, U is the reference potential, $C = C(x)$ is the so-called *doping profile*, τ_n, τ_p are the carrier life times, and

$$n_d = N_c e^{(E_t - E_c)/k_B T}, \quad p_d = N_v e^{(E_v - E_t)/k_B T}, \quad n_i = \sqrt{n_d p_d}.$$

The quantities N_c, N_v are the carrier effective densities of states, n_i is the intrinsic density, and E_t is the so-called *trap energy level*. The function $R(n, p)$ in (3) is called *Shockley-Read-Hall recombination-generation term*.

- (i) A thermal equilibrium state is a time-independent solution (n, p, V) to (1)–(4) with zero current flow, i.e.

$$\partial_t n = \partial_t p = 0, \quad J_n = J_p = 0.$$

For a generic equilibrium state, write n, p as functions only of the intrinsic density n_i , the a -dimensional potential V/U , and an integration constant. Write the equation satisfied by V .

- (ii) Consider (1)–(4) in a bounded domain Ω with the following Dirichlet boundary conditions:

$$n = n_D, \quad p = p_D, \quad V = V_D \quad \text{on } \partial\Omega, \quad (5)$$

where n_D, p_D, V_D are given functions. Under the assumptions that the densities are in equilibrium (recall point (i)) on $\partial\Omega$ and the total space charge $n_D - p_D - C$ vanish on $\partial\Omega$, write $n_D, p_D, V/U$ are functions only of C and n_i .

Exercise 2. Use the hints written at p. 21 of the lecture notes (about the Boltzmann equations) to derive the hydrodynamic model:

$$\partial_t n + \operatorname{div}(n\bar{v}) = 0, \quad (6)$$

$$\partial_t(n\bar{v}) + \operatorname{div}(n\bar{v} \otimes \bar{v}) + \frac{k_B}{m} \nabla(nT_e) + \frac{nqE}{m} = \int_{\mathbb{R}^3} vQ(f)dv, \quad (7)$$

$$\partial_t \mathcal{E} + \operatorname{div} \left(n\bar{v} \left(\frac{5}{2} k_B T_e + \frac{m|\bar{v}|^2}{2} \right) \right) + qn\bar{v} \cdot E = \int_{\mathbb{R}^3} \frac{m|v|^2}{2} Q(f)dv, \quad (8)$$

where n is the concentration, \bar{v} is the velocity, $\mathcal{E} = n \left(\frac{m|\bar{v}|^2}{2} + \frac{3}{2}k_B T_e \right)$ is the energy density, T_e is the effective temperature.

Exercise 3. The so-called *low-density* collision operator $Q_\ell(f)$ is defined as

$$Q_\ell(f)(v) = \frac{1}{\tau(x)} \left(\mathcal{M} \int_{\mathbb{R}^3} f(v') dv' - f(v) \right),$$

where $\tau(x) > 0$ is the relaxation time, \mathcal{M} is the “global Maxwellian”

$$\mathcal{M}(v) = \left(\frac{m}{2\pi k_B T_L} \right)^{3/2} e^{-m|v|^2/2k_B T_L},$$

and $T_L > 0$ is the (constant) lattice temperature.

- (i) Rewrite (6)–(8) with $Q = Q_\ell$.
- (ii) Consider now the system of equations:

$$\partial_t n + \operatorname{div}(n\bar{v}) = 0, \tag{9}$$

$$\partial_t(n\bar{v}) + \operatorname{div}(n\bar{v} \otimes \bar{v}) + \frac{k_B}{m} \nabla(nT_e) + \frac{nqE}{m} = -\frac{n\bar{v}}{\tau_p}, \tag{10}$$

$$\partial_t \mathcal{E} + \operatorname{div} \left(n\bar{v} \left(\frac{5}{2}k_B T_e + \frac{m|\bar{v}|^2}{2} \right) - \kappa n T_e \nabla T_e \right) + qn\bar{v} \cdot E = -\frac{n}{\tau_e} \left(\frac{\mathcal{E}}{n} - \frac{3}{2}k_B T_L \right), \tag{11}$$

where τ_p, τ_e are (x –dependent) relaxation times for momentum and energy, respectively, and $\kappa = \kappa(x) > 0$ is the heat conductivity. How do (9)–(11) relate with (6)–(8)? Rewrite (9)–(11) in terms of the a -dimensional variables $x_s, t_s, n_s, J_s, V_s, T_s$ defined as

$$x = \lambda x_s, \quad t = \tau t_s, \quad n = C_m n_s, \quad J = J_0 J_s, \quad V = UV_s, \quad T = T_L T_s,$$

where λ is the reference length (device diameter), C_m is the reference particle density, $U = k_B T_L / q$ is the reference potential, $J_0 = q C_m \lambda / \tau$ is the reference current density and $\tau = m^* \lambda / \tau_p k_B T_L$ is the reference time.

- (iii) Define the a -dimensional parameters $\alpha = \tau_p / \tau, \beta = \tau_e / \tau$. Take the limit $\alpha \rightarrow 0$ in the scaled version of (9)–(11); this limit system is called (scaled) *energy-transport model*. If you take the limit $\beta \rightarrow 0$ in the scaled energy-transport model what do you obtain?

The exercises will be reviewed in class on January 18th, 2017.