

# Elastic Curves in Riemannian Space

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# Agenda

Length of Curves

Elastic Energy

Willmore Surfaces

Bibliography

## Length of Curves

Let  $(M, g = \langle \cdot, \cdot \rangle)$  be a Riemannian manifold,  $I \in \{[a, b], \mathbb{S}^1\}$ ,  $\gamma : I \rightarrow M$  smooth.  $\dot{\gamma} = T\gamma(\frac{d}{dt})$  is the push-forward of the vector field  $\frac{d}{dt}$ , i.e. the tangent vector field in  $\mathcal{T}(\gamma)$ .

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### Definition

The length of  $\gamma$  is given by

$$\mathcal{L}(\gamma) := \int_I |\dot{\gamma}(t)| dt = \int_I g_{\gamma(t)}(\dot{\gamma}(t), \dot{\gamma}(t))^{\frac{1}{2}} dt.$$

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Further assume that  $\gamma$  is an immersion, i.e.  $|\dot{\gamma}| \neq 0$  on  $I$ ), and  $\gamma$  is parametrized s.t.  $|\dot{\gamma}| = 1$ . Since  $M$  is Riemannian there exists a unique connection  $\nabla$  yielding  $D_t$ , the covariant derivative along  $\gamma$ .

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### Definition

The curvature vector of  $\gamma$  is defined by  $D_t \dot{\gamma} \in \mathcal{T}(\gamma)$ .

## Extremize $\mathcal{L}$

Let  $W \in \mathcal{T}(\gamma)$  be a smooth vector field along  $\gamma$ .

The geodesic flow then yields the existence of a smooth variation  $\gamma_\varepsilon(\cdot) : I \rightarrow M$  such that

$$\left. \frac{d}{d\varepsilon} \gamma_\varepsilon \right|_{\varepsilon=0} = W \text{ and } \gamma_0 = \gamma.$$

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Thus for the first variation of  $\mathcal{L}$  we have by compatibility

$$\left. \frac{d}{d\varepsilon} \mathcal{L}(\gamma_\varepsilon) \right|_{\varepsilon=0} = [\langle W, \dot{\gamma} \rangle]_0^{\mathcal{L}(\gamma)} - 2 \int_0^{\mathcal{L}(\gamma)} \langle D_t \dot{\gamma}, W \rangle ds$$

where  $s \in [0, \mathcal{L}(\gamma)]$  denotes the arc length, i.e.  $ds = \|\dot{\gamma}\| dt$ .



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where  $s \in [0, \mathcal{L}(\gamma)]$  denotes the arc length, i.e.  $ds = \|\dot{\gamma}\| dt$ . Hence we see for fixed end points ( $W|_{\partial I} \equiv 0$ ):

### Proposition

*The curve  $\gamma$  is critical (i.e.  $\left. \frac{d}{d\varepsilon} \mathcal{L}(\gamma_\varepsilon) \right|_{\varepsilon=0} = 0 \forall W \in \mathcal{T}(\gamma)$ )  
 $\Leftrightarrow D_t \dot{\gamma} = 0$ , i.e. if  $\gamma$  is geodesic.*

## Elastic Energy

Let  $\gamma : I \rightarrow M$  be an immersion.

### Definition

The Elastic Energy (or Bending Energy) of  $\gamma$  is defined as

$$\mathcal{E}(\gamma) := \int_0^{\mathcal{L}(\gamma)} |D_t \dot{\gamma}|^2 ds$$

For any  $\lambda > 0$ , an Elastica (or Elastic Curve) is an immersion  $\gamma : I \rightarrow M$  which is critical for  $\mathcal{E}(\cdot) + \lambda \mathcal{L}(\cdot) =: \mathcal{E}^\lambda(\cdot)$ .

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### Definition

For a vector field  $X \in \mathcal{T}(\gamma)$  we denote the normal component of  $D_t \dot{\gamma}$  by

$$\nabla^\perp X := D_t X - \langle D_t X, \dot{\gamma} \rangle \dot{\gamma}.$$

With help of this abbreviation we can now formulate the first variation of  $\mathcal{E}^\lambda$  in the next proposition.

## The First Variation

### Proposition

Let  $\gamma_\varepsilon$  be a smooth variation of  $\gamma$  with variation field  $W$ . Then the following equation holds:

$$\left. \frac{d}{d\varepsilon} \mathcal{E}^\lambda(\gamma_\varepsilon) \right|_{\varepsilon=0}$$

where  $R$  is the Riemannian curvature endomorphism.

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where  $R$  is the Riemannian curvature endomorphism. For closed curves  $\gamma : \mathbb{S}^1 \rightarrow M$  the boundary term vanishes and thus the Euler-Lagrange Equation is given by

$$2R(D_t \dot{\gamma}, T)T + 2(\nabla^\perp)^2 D_t \dot{\gamma} + (|D_t \dot{\gamma}|^2 - \lambda) D_t \dot{\gamma} = 0 \text{ on } \gamma. \quad (1)$$

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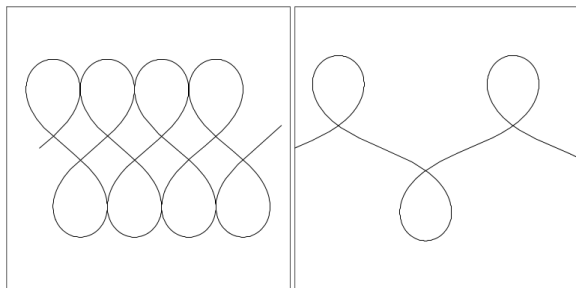
$$\frac{d}{d\varepsilon} \mathcal{E}^\lambda(\gamma_\varepsilon) \Big|_{\varepsilon=0} = \int_0^{\mathcal{L}(\gamma)} \langle 2R(D_t \dot{\gamma}, \dot{\gamma}) \dot{\gamma} + 2(\nabla^\perp)^2 D_t \dot{\gamma} + (|D_t \dot{\gamma}|^2 - \lambda) D_t \dot{\gamma}, W \rangle ds \\ + \left[ 2 \langle \nabla^\perp W, D_t \dot{\gamma} \rangle - \langle 2\nabla^\perp D_t \dot{\gamma} + (|D_t \dot{\gamma}|^2 - \lambda) \dot{\gamma}, W \rangle \right]_0^{\mathcal{L}(\gamma)}.$$

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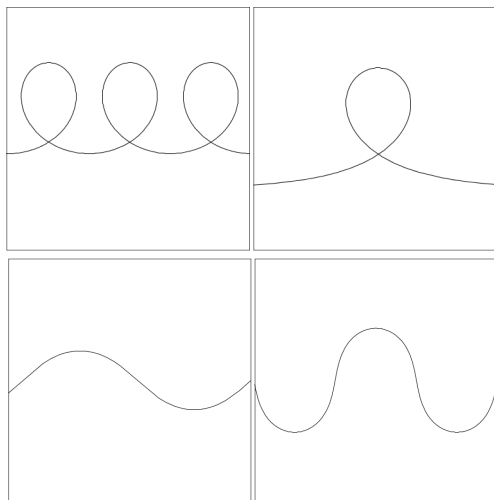
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For space forms, where the sectional curvature is constant to a real number  $K_0$  the Euler-Lagrange Equation (1) simplifies to

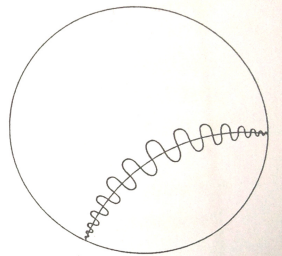
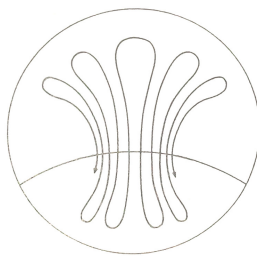
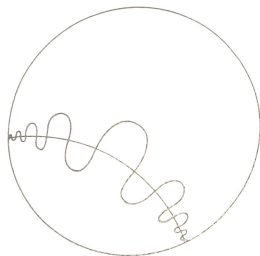
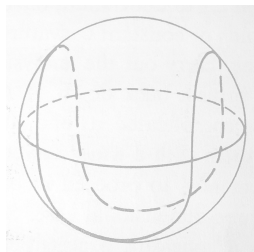
$$2(\nabla^\perp)^2 D_t \dot{\gamma} + (|D_t \dot{\gamma}|^2 - \lambda + 2K_0) D_t \dot{\gamma} = 0 \text{ on } \gamma.$$



(Taken from [1])



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(Taken from [2])

## Willmore Surfaces

Let  $\iota : \Sigma^2 \hookrightarrow \mathbb{R}^3$  be an immersion,  $\Sigma^2$  a two-dimensional compact manifold.

### Definition

$\Sigma$  is called Willmore Surface if it is critical for

$$\mathcal{W}(\Sigma) = \int_{\Sigma} H^2 \, d\text{Vol}_g$$

where  $g = \iota^* g_{\text{Eucl}}$  is the induced metric and  $H = \frac{1}{2} \text{tr}(S) = \frac{1}{2}(\kappa_1 + \kappa_2)$  is the mean curvature of  $\Sigma$ .

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### Proposition

The Euler-Lagrange equation for  $\mathcal{W}$  is given by

$$\Delta H + 2H(H^2 - K) \text{ on } \Sigma \quad (2)$$



## Proposition (Langer, Singer '84 [3])

Let  $\gamma$  be a closed curve in the hyperbolic half-plane. Then for the torus of revolution  $\Gamma$  of  $\gamma$  in  $\mathbb{R}^3$  we have  $\mathcal{W}(\Gamma) = \frac{\pi}{2} \mathcal{E}_{\mathbb{H}}(\gamma)$ .

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### Proposition (Dall'Acqua, Deckelnick, Grunau '08 [4])

For a graph  $u : [-1, 1] \rightarrow \mathbb{R}^+$  smooth consider the surface of revolution given by

$$\Gamma := \left\{ \begin{pmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} u(x) \\ 0 \\ x \end{pmatrix}, \phi \in \mathbb{R}, x \in [-1, 1] \right\}$$

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Then

$$\mathcal{W}(\Gamma) = \frac{\pi}{2} \mathcal{E}_{\mathbb{H}}^0 \left( \begin{pmatrix} u(\cdot) \\ \cdot \end{pmatrix} \right) - 2\pi \left[ \frac{\dot{u}(t)}{\sqrt{1 + \dot{u}^2(t)}} \right]_{-1}^1$$

where  $\mathcal{E}_{\mathbb{H}}$  denotes the elastic energy of the curve  $[-1, 1] \ni t \mapsto (u(t), t) \in \mathbb{H}$  in the hyperbolic half-plane.

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